

VDEs and Resistive Wall Instabilities with M3D-C1

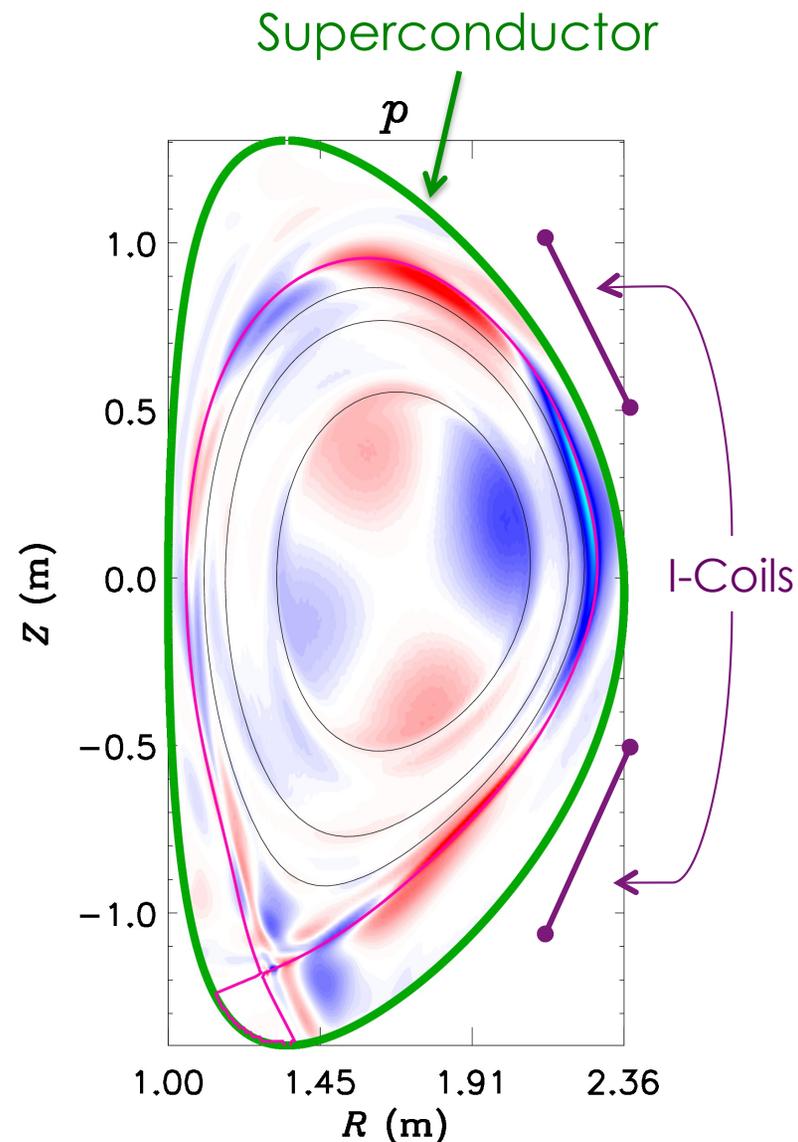
by
N.M. Ferraro

Presented at the
Theory and Simulation of Disruptions Workshop

Princeton, NJ
July 10, 2014

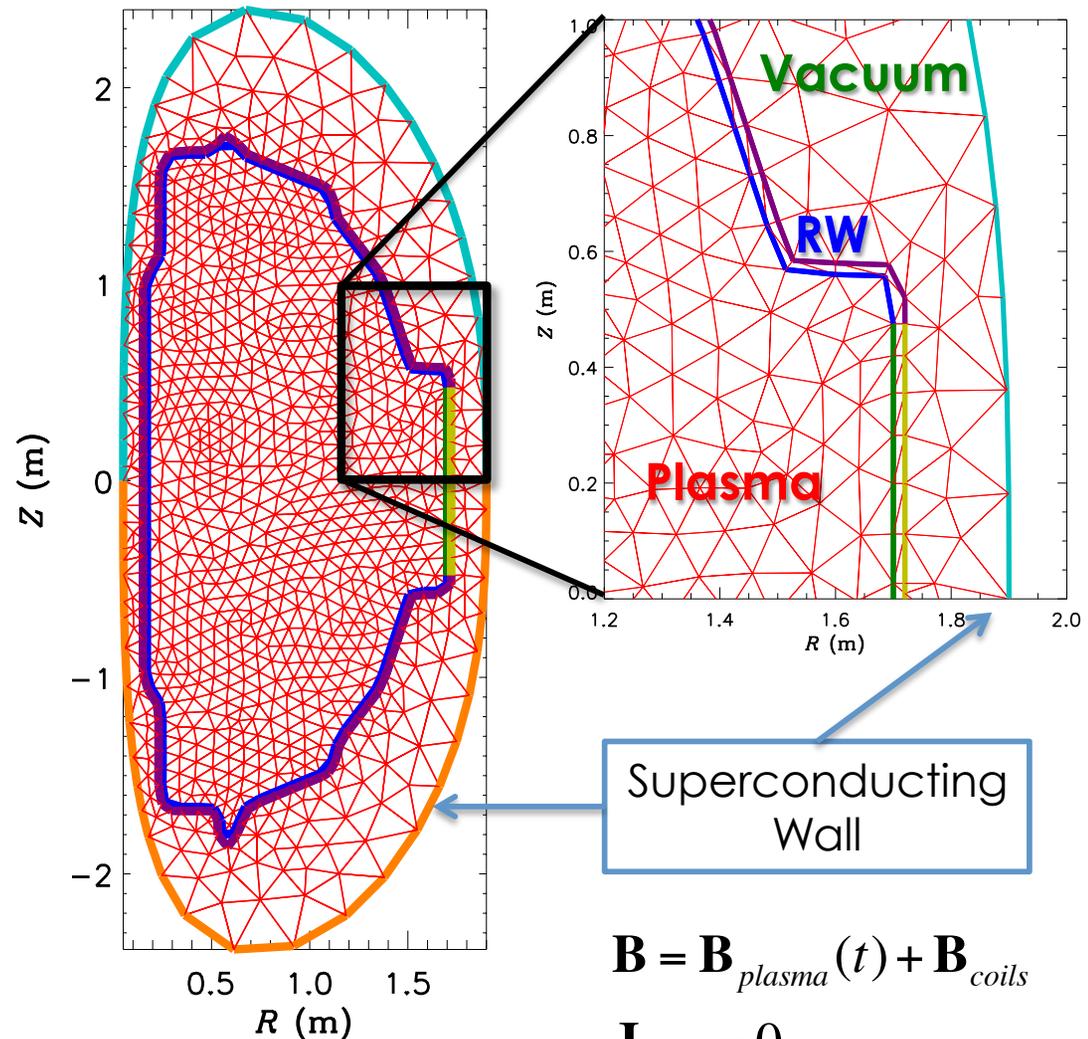
Previous M3D-C1 Calculations Required Conducting Boundary Within PF Coils

- **Affects stability & plasma response**
 - Especially $n = 0, n = 1$
- **Shields magnetic probes from plasma response**
 - Plasma response outside of conductor is zero
- **Implementing resistive wall boundary conditions was challenging**
 - All boundary nodes become coupled; hurts parallel scalability
 - Extant RW codes are spectral



New Resistive Wall Capability In M3D-C1

- **3 regions inside domain:**
 - Vacuum ($\mathbf{J} = 0$)
 - RW ($\mathbf{E} = \eta_w \mathbf{J}$)
 - Plasma (Extended MHD)
- **Boundary conditions:**
 - v, p, n set at inner wall
 - \mathbf{B} set at outer (superconducting) wall
- **There are no boundary conditions on \mathbf{B} or \mathbf{J} at the resistive wall**
 - Current can flow into and through the wall



$$\mathbf{B} = \mathbf{B}_{plasma}(t) + \mathbf{B}_{coils}$$

$$\mathbf{J}_{coils} = 0$$

Advantages and Disadvantages of Including Resistive Wall In Domain

- **Advantages:**

- Computation is more scalable than using RW BCs for implicit step
 - RW BCs couple all finite elements touching the boundary
- Can add time/space dependent physical attributes of wall
 - Resistivity, temperature
- Can allow current to flow into and out of wall

- **Disadvantages:**

- Bigger matrices
 - But non-MHD regions do not make matrices more poorly conditioned
- Still need a conducting boundary somewhere
 - This could be a problem in STs like NSTX-U

Two-Fluid Model is Implemented in “Plasma” Region

$$\frac{\partial n}{\partial t} + \nabla \cdot (n_i \mathbf{v}) = 0$$

$$n_i m_i \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \cdot \Pi_i$$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p + \Gamma p \nabla \cdot \mathbf{v} = -\frac{1}{n_e e} \mathbf{J} \cdot \left(\Gamma p_e \frac{\nabla n_e}{n_e} - \nabla p_e \right) - (\Gamma - 1) \nabla \cdot \mathbf{q}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

$$\mathbf{E} = -\mathbf{v} \times \mathbf{B} + \eta \mathbf{J} + \frac{1}{n_e e} (\mathbf{J} \times \mathbf{B} - \nabla p_e)$$

$$\Pi_i = -\mu \left[\nabla \mathbf{v} + (\nabla \mathbf{v})^T \right] + \Pi_i^{gv} + \Pi_i^{\parallel}$$

$$\mathbf{q} = -\kappa \nabla T_i - \kappa_{\parallel} \mathbf{b} \mathbf{b} \cdot \nabla T_e$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

$$\Gamma = 5/3$$

$$n_e = \sum_i Z_i n_i$$

- Ion viscosity model optionally includes Braginskii gyroviscosity, parallel viscosity (poloidal flow damping)
- Open field line region of “plasma” region is treated as low-temperature, low-density plasma

Two-Fluid Model is Implemented in “Plasma” Region

$$\frac{\partial n}{\partial t} + \nabla \cdot (n_i \mathbf{v}) = 0$$

$$n_i m_i \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \cdot \Pi_i$$

~~$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p + \Gamma p \nabla \cdot \mathbf{v} = -\frac{1}{n_e e} \mathbf{J} \cdot \left(\Gamma p_e \frac{\nabla n_e}{n_e} - \nabla p_e \right) - (\Gamma - 1) \nabla \cdot \mathbf{q}$$~~

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

~~$$\mathbf{E} = -\mathbf{v} \times \mathbf{B} + \eta \mathbf{J} + \frac{1}{n_e e} (\mathbf{J} \times \mathbf{B} - \nabla p_e)$$~~

~~$$\Pi_i = -\mu \left[\nabla \mathbf{v} + (\nabla \mathbf{v})^T \right] + \Pi_i^{sv} + \Pi_i^{\parallel}$$~~

$$\mathbf{q} = -\kappa \nabla T_i - \kappa_{\parallel} \mathbf{b} \mathbf{b} \cdot \nabla T_e$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

$$\Gamma = 5/3$$

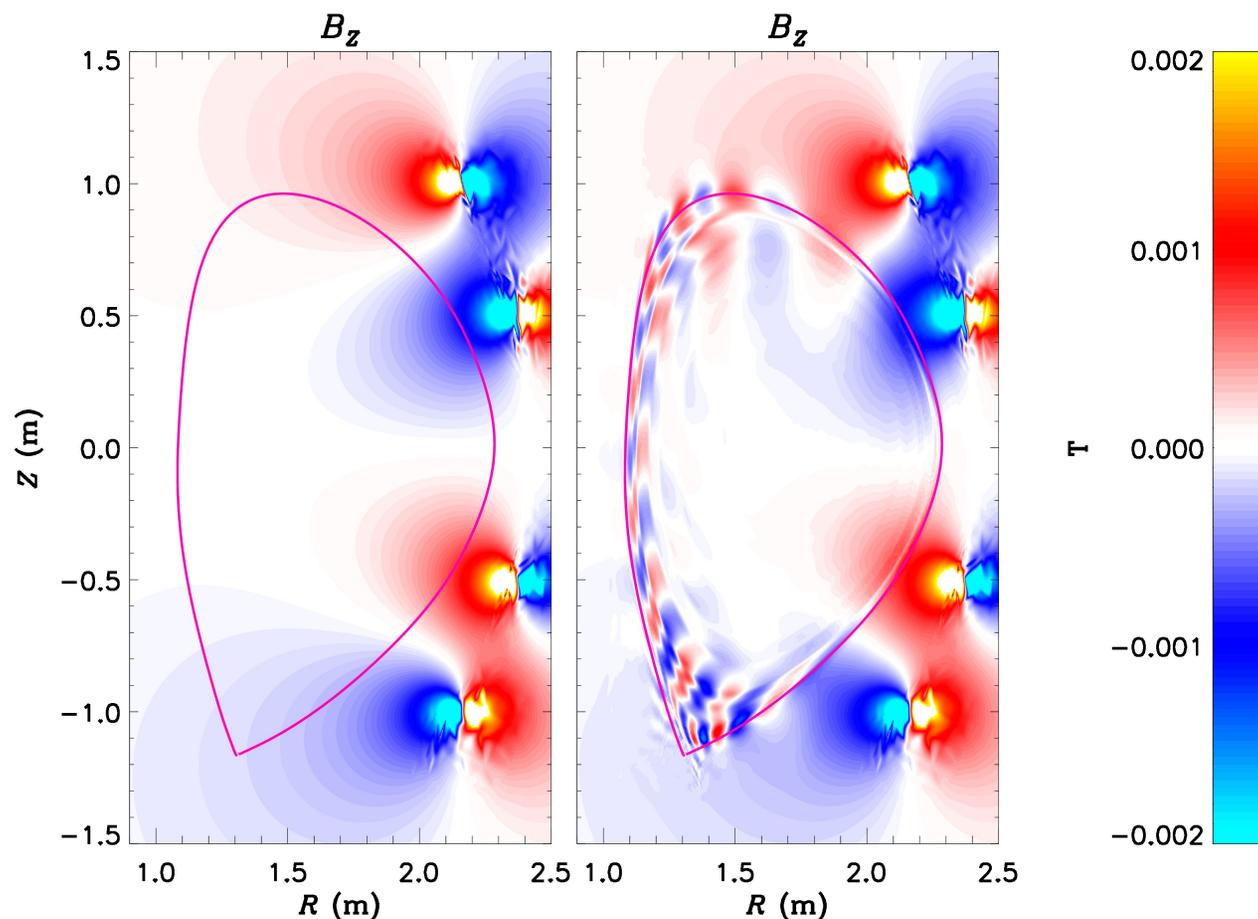
$$n_e = \sum_i Z_i n_i$$

- Ion viscosity model optionally includes Braginskii gyroviscosity, parallel viscosity (poloidal flow damping)
- Open field line region of “plasma” region is treated as low-temperature, low-density plasma
- VDE calculations here use a single-fluid model

“Free Boundary” 3D Response

Resistive Wall Capability Allows “Free-Boundary” 3D Response Calculations

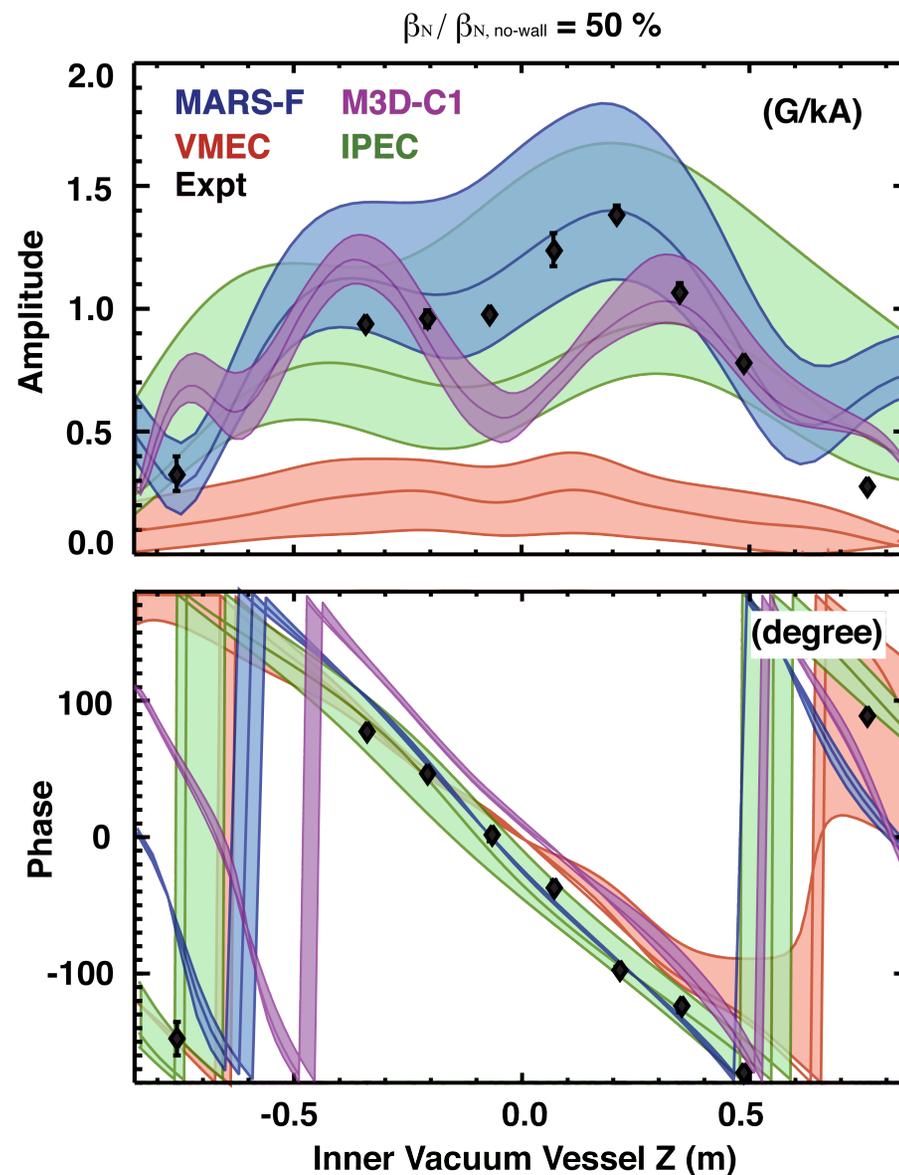
- (Technically fixed-boundary, but now conducting wall is very far from plasma and outside PF coils)
- In zero-frequency response, there are no eddy currents
 - Resistive wall can still play a role: currents can flow through wall
- Free-boundary response allows direct comparison with new MP data



NM Ferraro/TSD/July 2014

Preliminary Free-Boundary Calculations Show Encouraging Agreement With Experimental Data

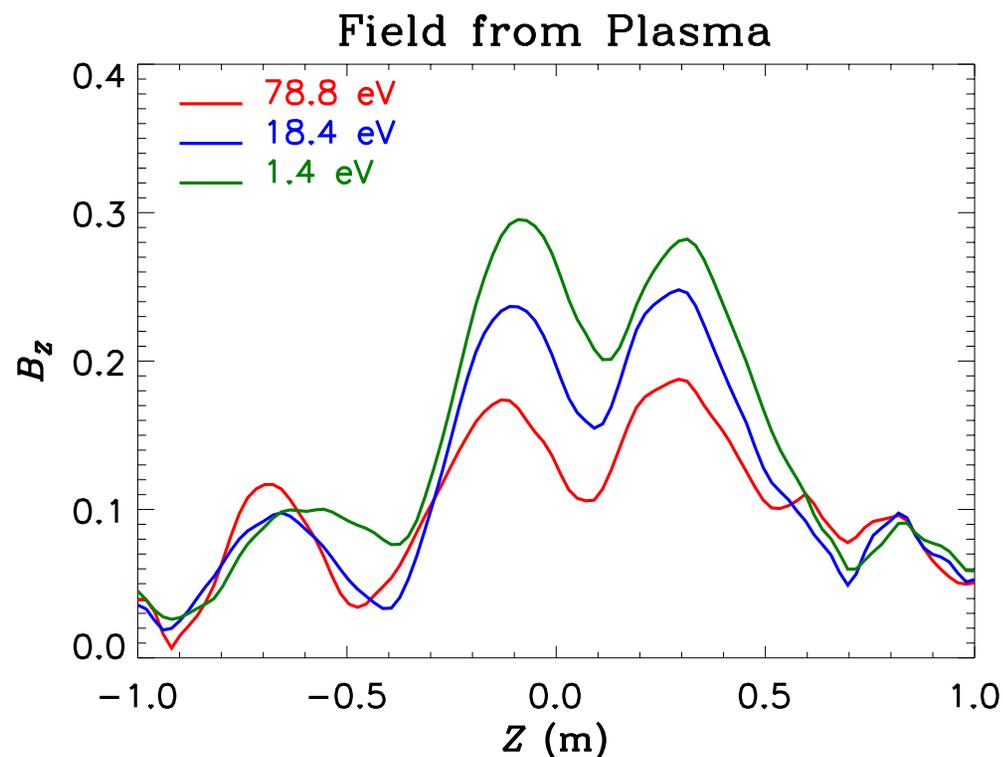
- **Free-Boundary 3D response calculated with several codes**
 - IPEC, MARS-F: linear, ideal
 - M3D-C1: linear, two-fluid, resistive
 - VMEC: nonlinear, ideal
- **Calculated values are in decent agreement with measurements**
- **Different codes show different sensitivities to bootstrap current**
 - M3D-C1 seems least sensitive, probably because there is no “ q_{edge} ”



Courtesy J. King

Inclusion of Open Field Line Region Introduces Additional “Free Parameters” in M3D-C1 Model

- Because M3D-C1 models open field-line region as a plasma, the parameters of this region can affect the response
- Magnetic probe response is especially sensitive to the resistivity of this region
- Density is much less important
- Validation is ongoing, and is helping to refine the models

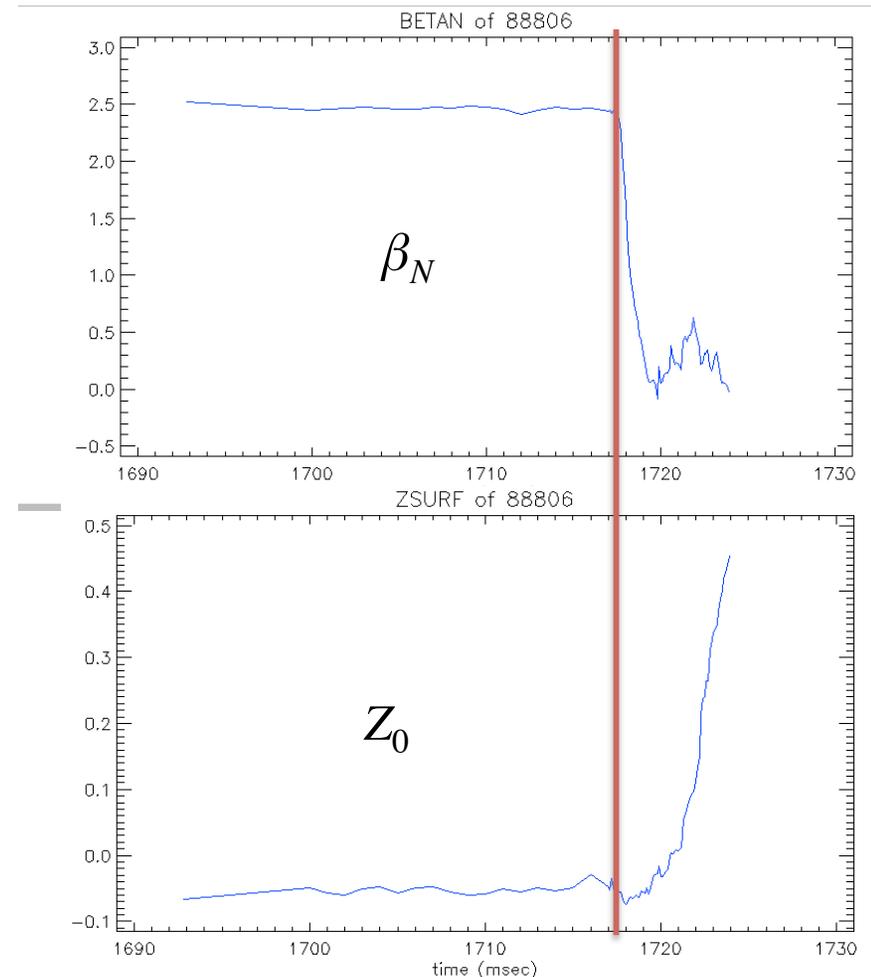
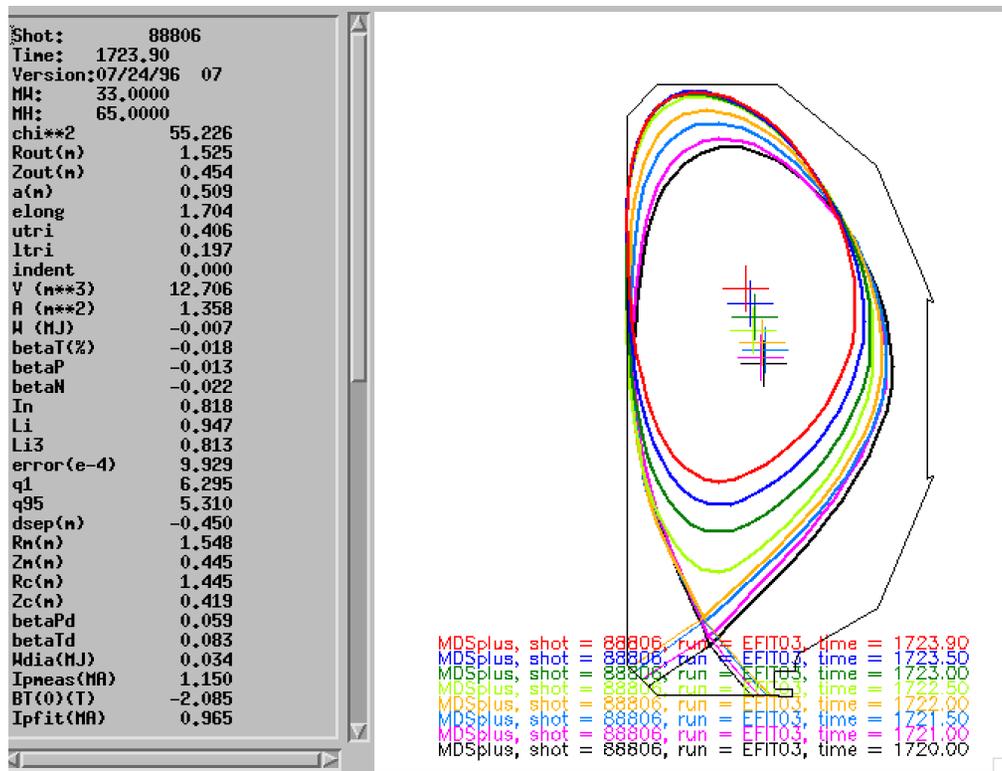


NM Ferraro/TSD/July 2014

Vertical Displacement Events

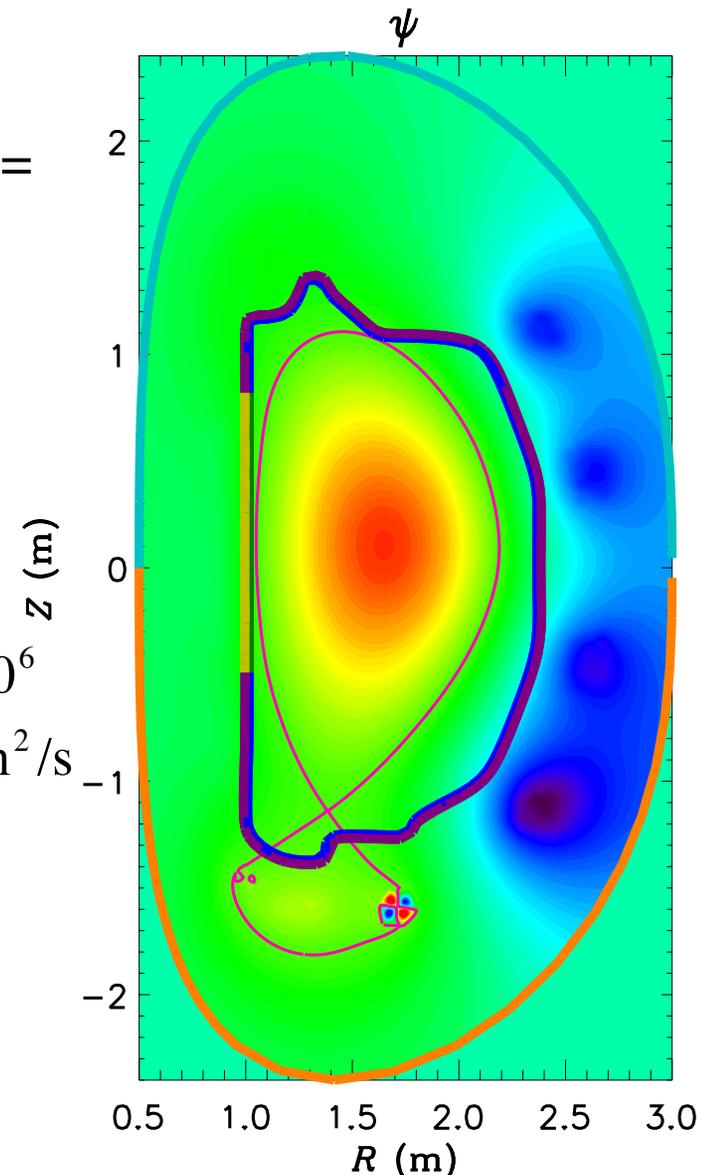
Nonlinear Calculation Recovers $n = 0$ Instability In DIII-D VDE Discharge

- DIII-D discharge 088806 disrupted due to gas injection
 - Vertical stability was lost shortly after thermal quench
 - Timescale ~ 3 ms



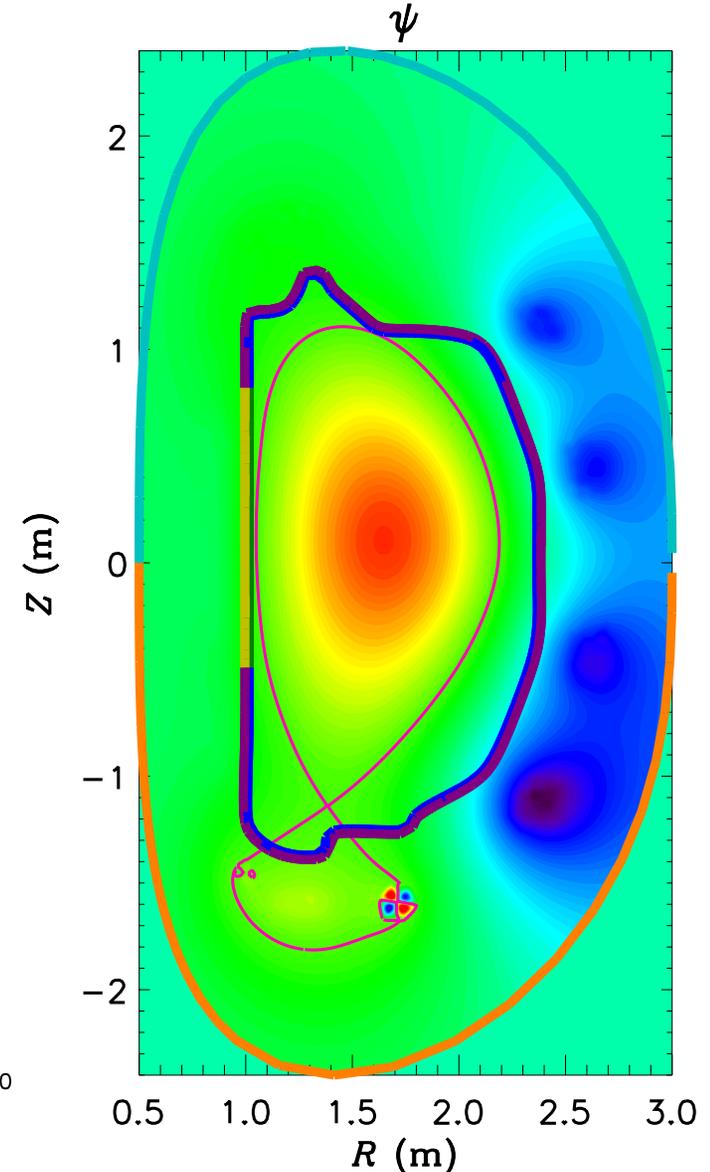
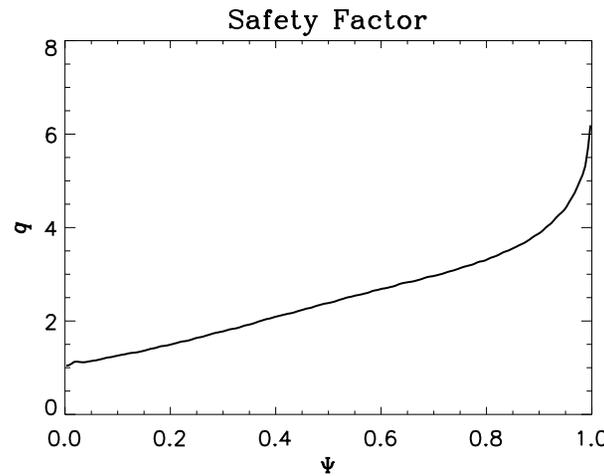
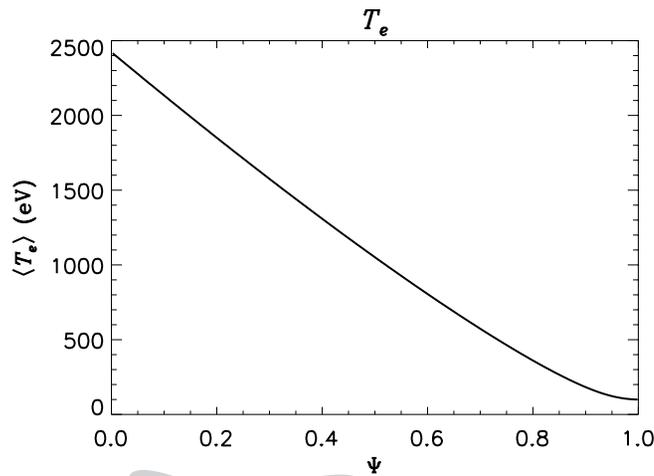
Nonlinear Calculation Initialized From EFIT Reconstruction

- **M3D-C1 was initialized using the reconstructed equilibrium just before TQ ($t = 1720$ ms)**
 - Equilibrium is re-solved on M3D-C1 grid
- **Nonlinear $n = 0$ calculation uses fairly realistic plasma parameters**
 - Spitzer resistivity: $S_0 \approx 6.8 \times 10^{-7}$
 - Anisotropic thermal conductivity: $\chi_{\parallel} / \chi_{\perp} = 10^6$
 - Anomalous perp. transport: $100 < \chi_{\perp} < 800 \text{ m}^2/\text{s}$
- **RW approximates first wall, not vacuum vessel here; using “modern” first wall, different from old experiment**



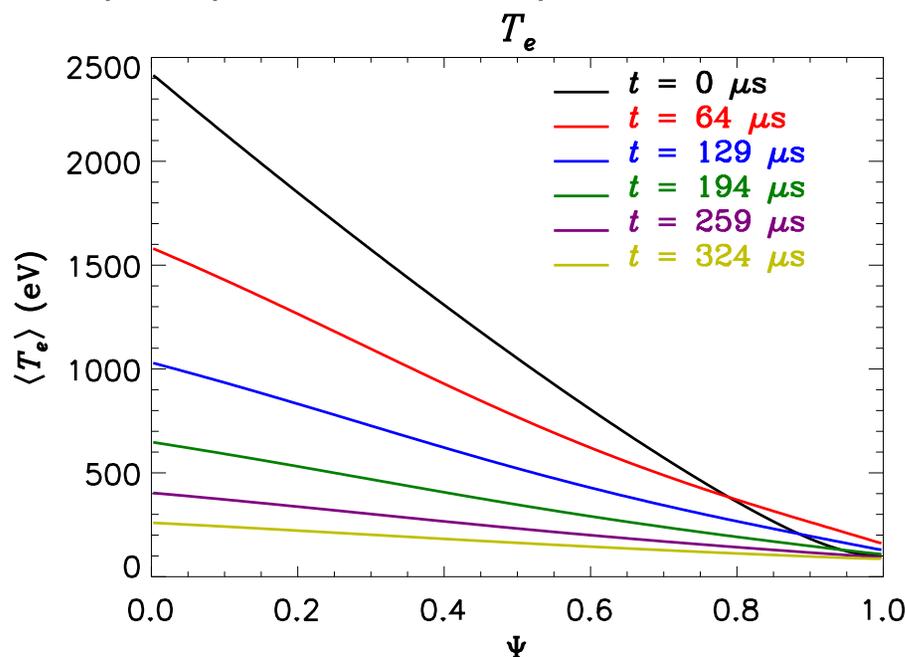
These Calculations Are A “First Try”; Not Suitable For Quantitative Validation

- **Simulations done at low resolution**
 - 5059 elements, ~320k DOFs
- $T_{SOL} \approx 100 \text{ eV} \rightarrow \eta_{SOL} \approx 1.6 \times 10^{-6} \Omega \text{ m}$
- **Single-Fluid, no sources**
- **Wall is uniform thickness (2 cm), resistivity**



Simulations Include Thermal Quench Stage

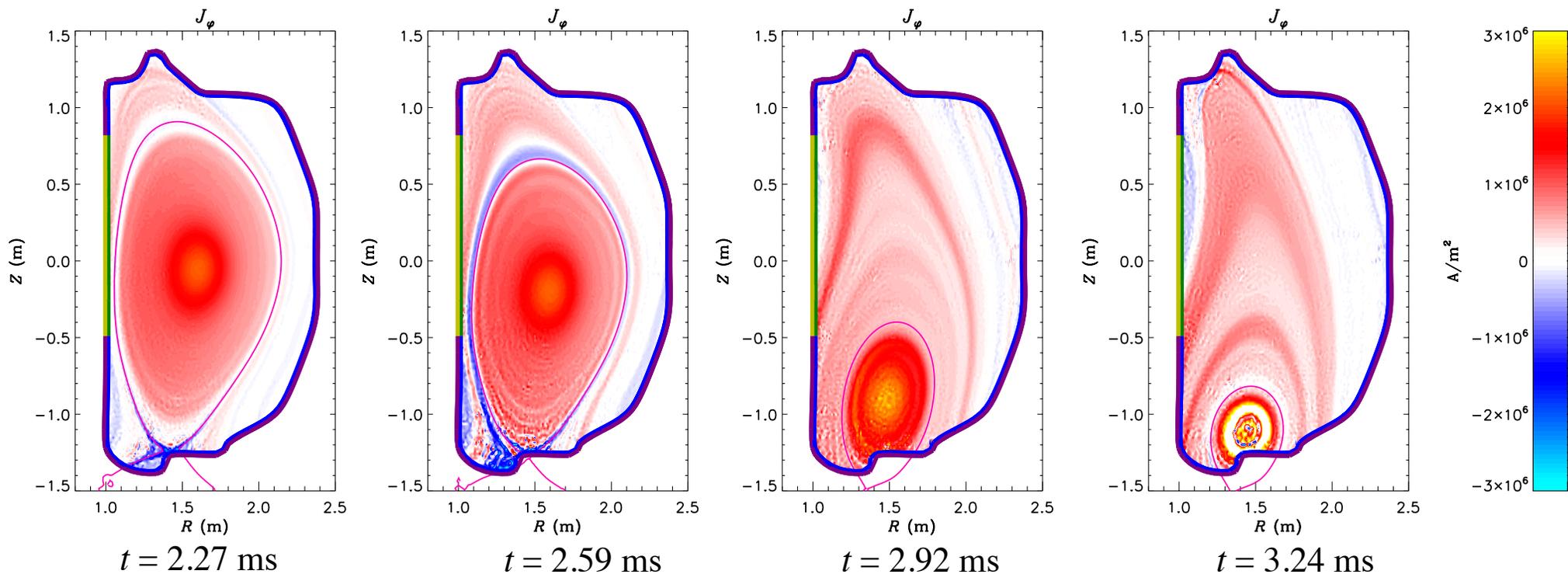
- A thermal collapse happens on $\sim 100 \mu\text{s}$ timescale, due to large perpendicular thermal conductivity
 - Not caused by any MHD activity or convective transport



- At some point during the TQ, the plasma becomes vertically unstable

Calculation Shows Vertical Displacement Into Lower Divertor

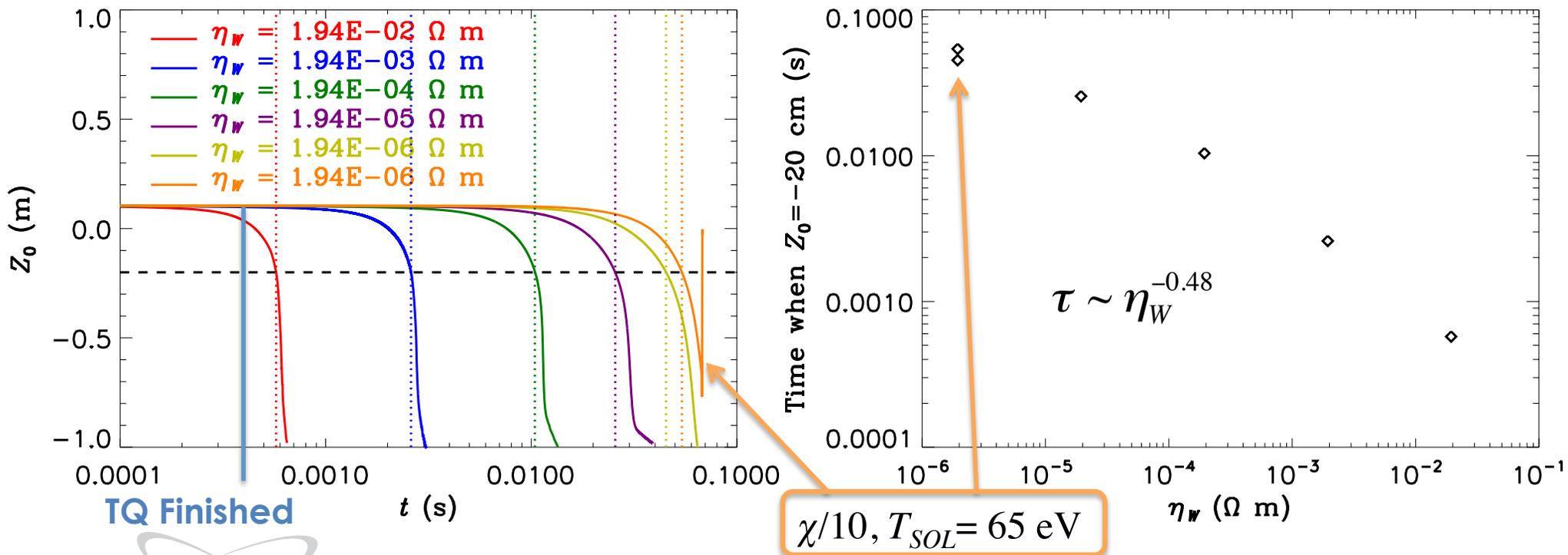
- Both **co- I_p** and **counter- I_p** currents are seen in the open field-line region
- **Plasma always moves to lower divertor, unlike in experiment**
 - Maybe due to different wall configuration?



Timescale of VDE Scales Inversely with $(\eta_w)^{1/2}$

- Given wall thickness $\delta = 2$ cm and a poloidal scale length $d = 50$ cm, resistive wall diffusion times range from ~ 6.5 ms to ~ 0.65 μ s
- VDE timescale is longer than resistive wall time
 - Doesn't seem strongly affected by T_{SOL} ; need more cases

$$\tau_w = \frac{\mu_0 d \delta}{\eta_w}$$

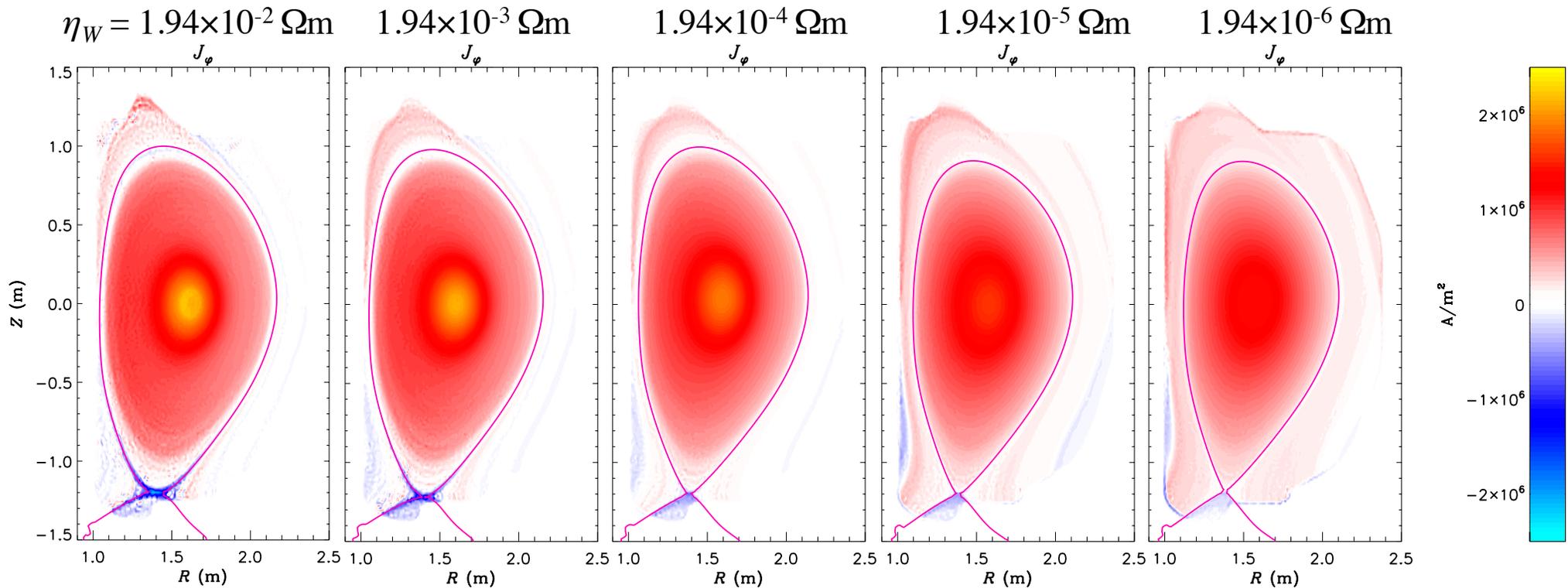


TQ Finished



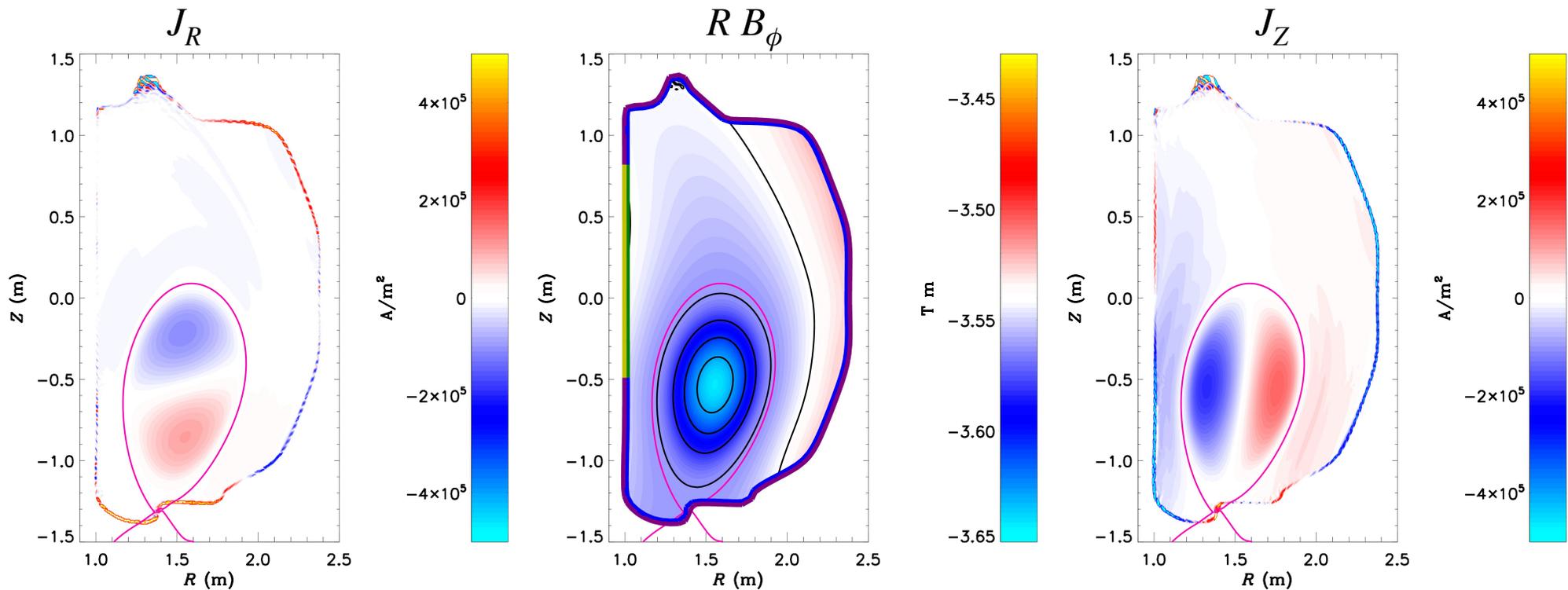
Currents in Wall and Open Field-Line Region Change with η_W

- At early stage of VDE, currents in the wall are stronger at lower η_W
- **Counter- I_p** currents are significantly stronger at higher η_W



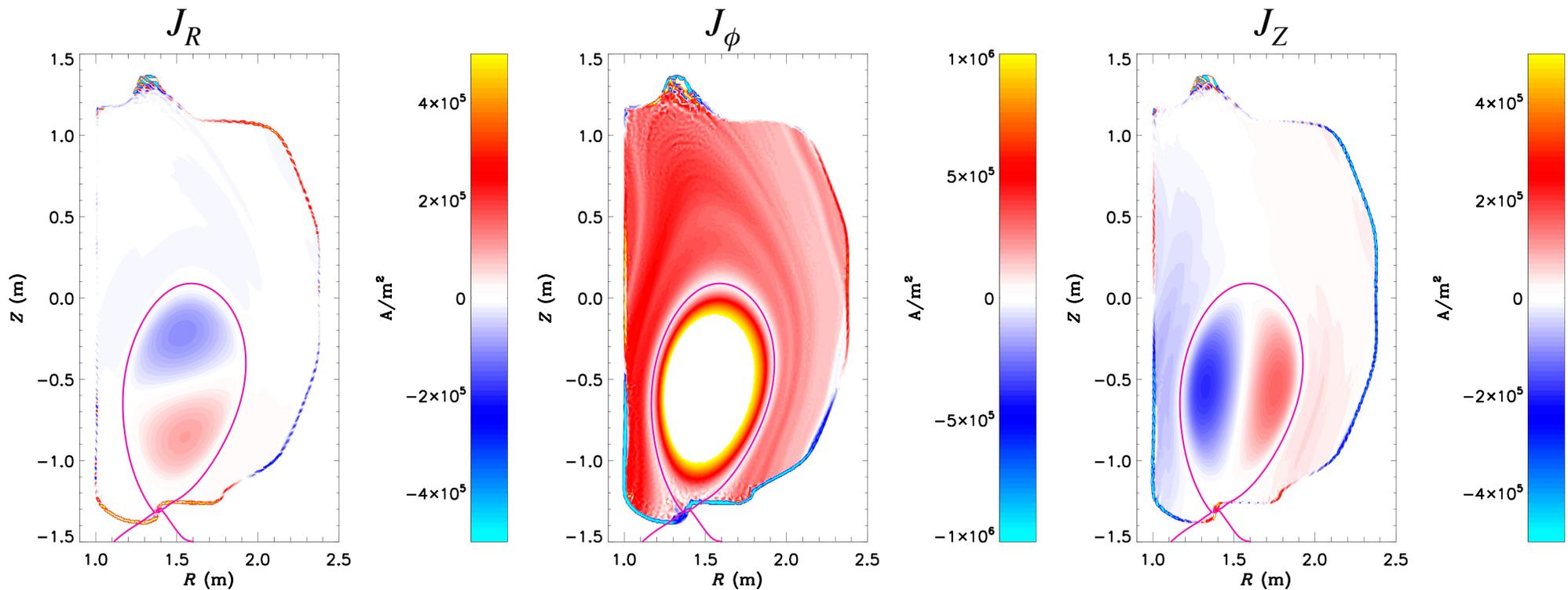
Wall Currents are Mostly Inductive

- **Poloidal currents are present in the open field-line region**
 - Gradients in $R B_\phi$ imply poloidal currents
 - Current flows from plasma to wall to ensure $\nabla \cdot \mathbf{J} = 0$
- **Poloidal wall currents are consistent with excluding toroidal flux**



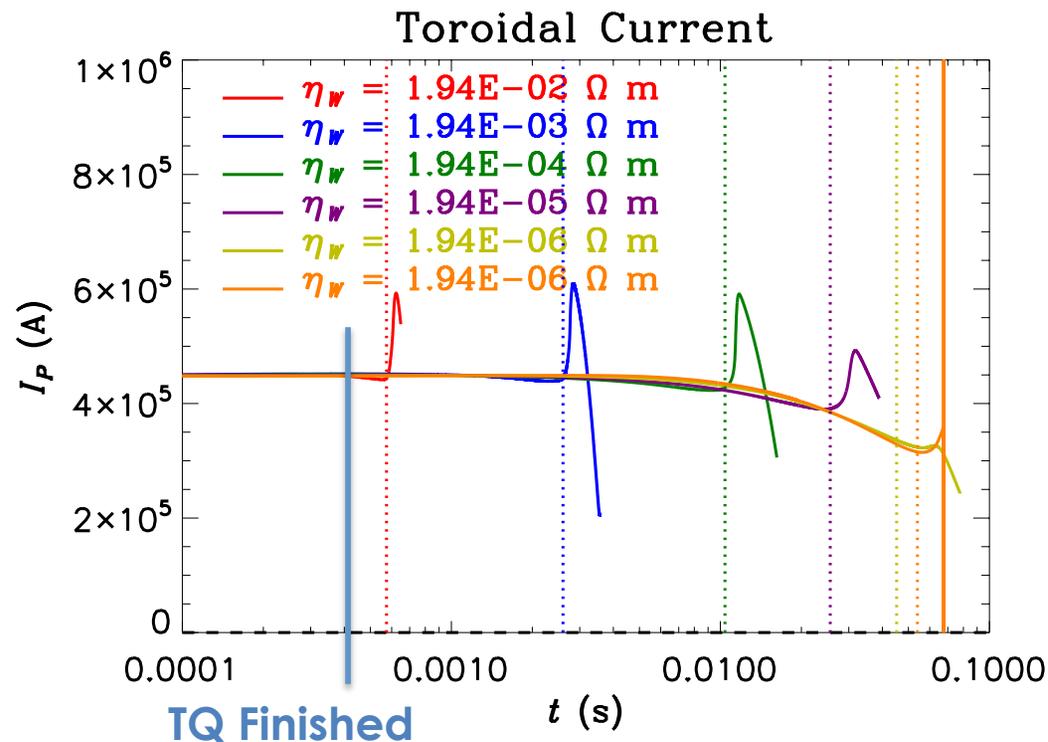
Wall Currents are Mostly Inductive

- **Toroidal currents are also present in the open field-line region**
 - Magnitude may be an artifact of high T_e in the open field-line region
- **Toroidal wall currents are consistent with excluding poloidal flux**



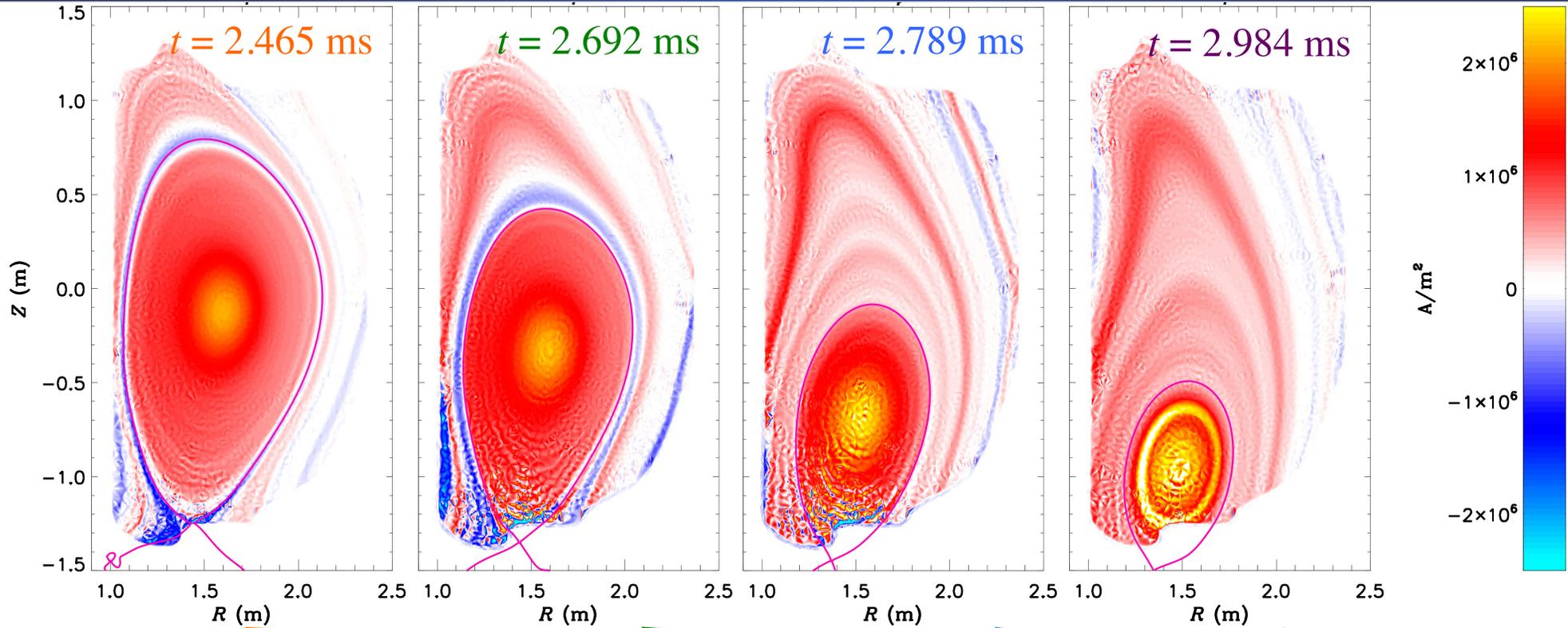
Current Spikes Observed Before Current Quench; Associated with Vertical Motion of Plasma

- Current spike onset is correlated with vertical motion of plasma, unlike TQ



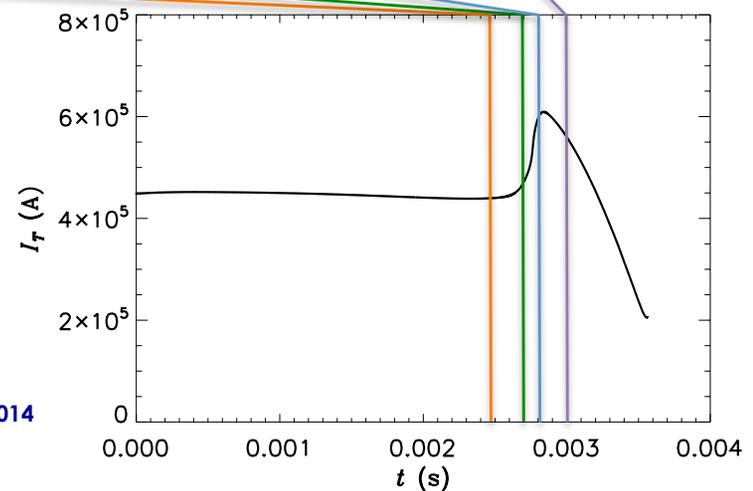
- “ I_P ” here only includes all toroidal current in the plasma region, but not in the resistive wall

Current Spike is Associated With Loss of Counter- I_p Current In Open Field-Line Region

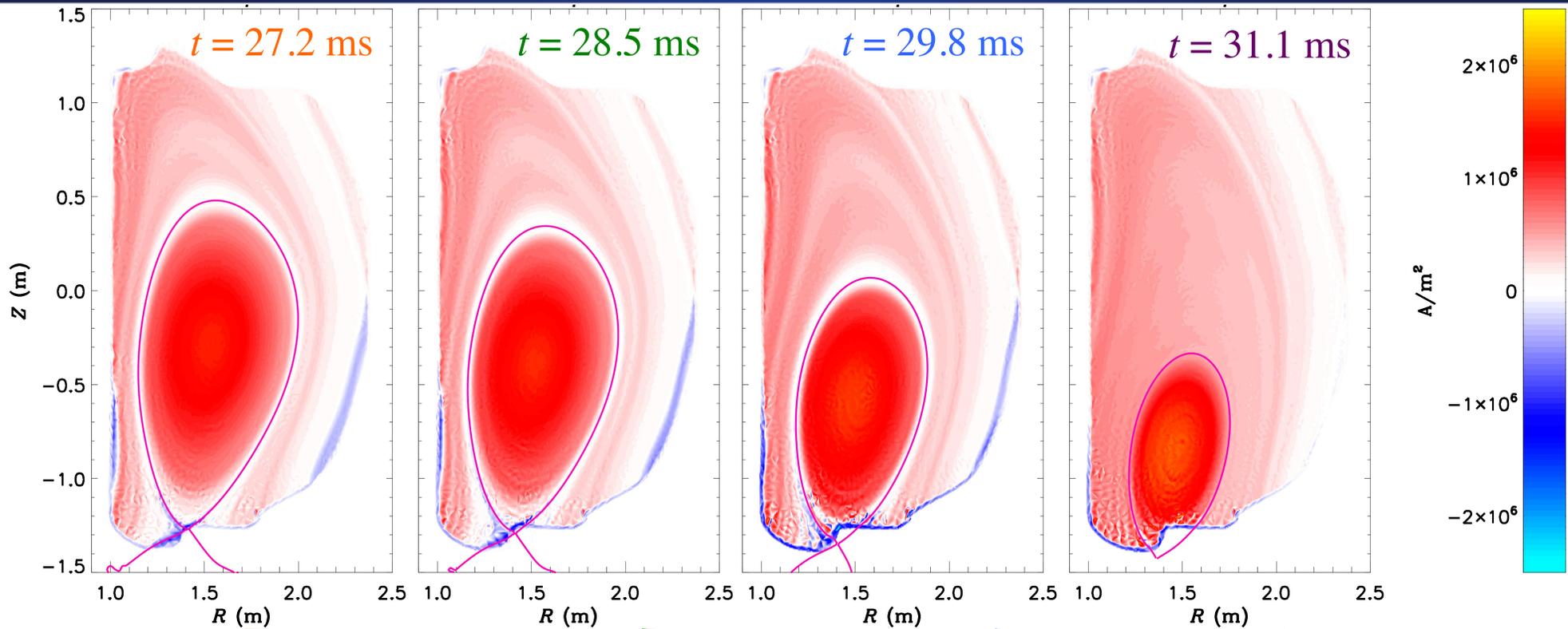


$$\eta_W = 1.94 \times 10^{-3} \Omega \text{ m}$$

- Plasma undergoes rapid contraction during current spike

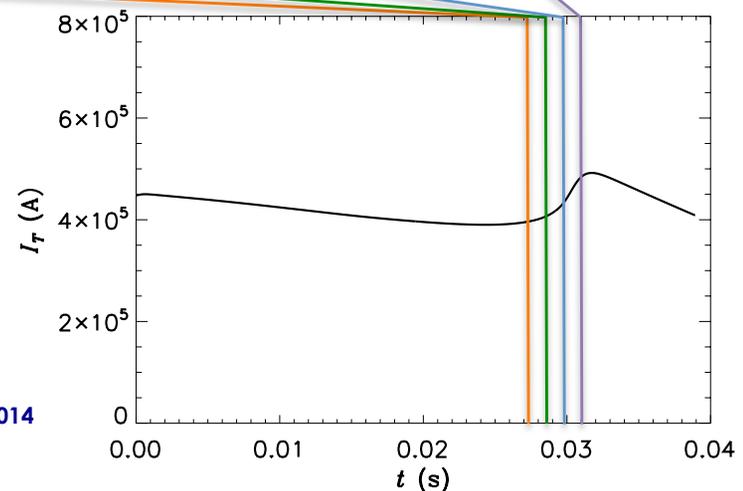


Current Spike is Associated With Loss of Counter- I_p Current In Open Field-Line Region



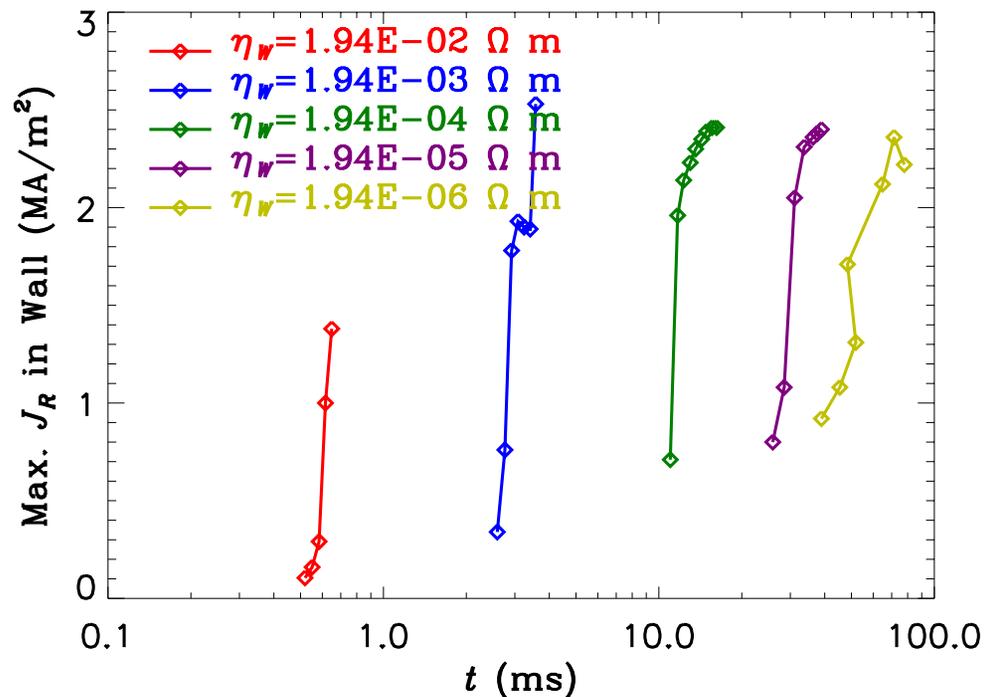
$$\eta_W = 1.94 \times 10^{-6} \Omega \text{ m}$$

- Plasma undergoes rapid contraction during current spike



Max Poloidal Current in Wall Depends Weakly on η_W

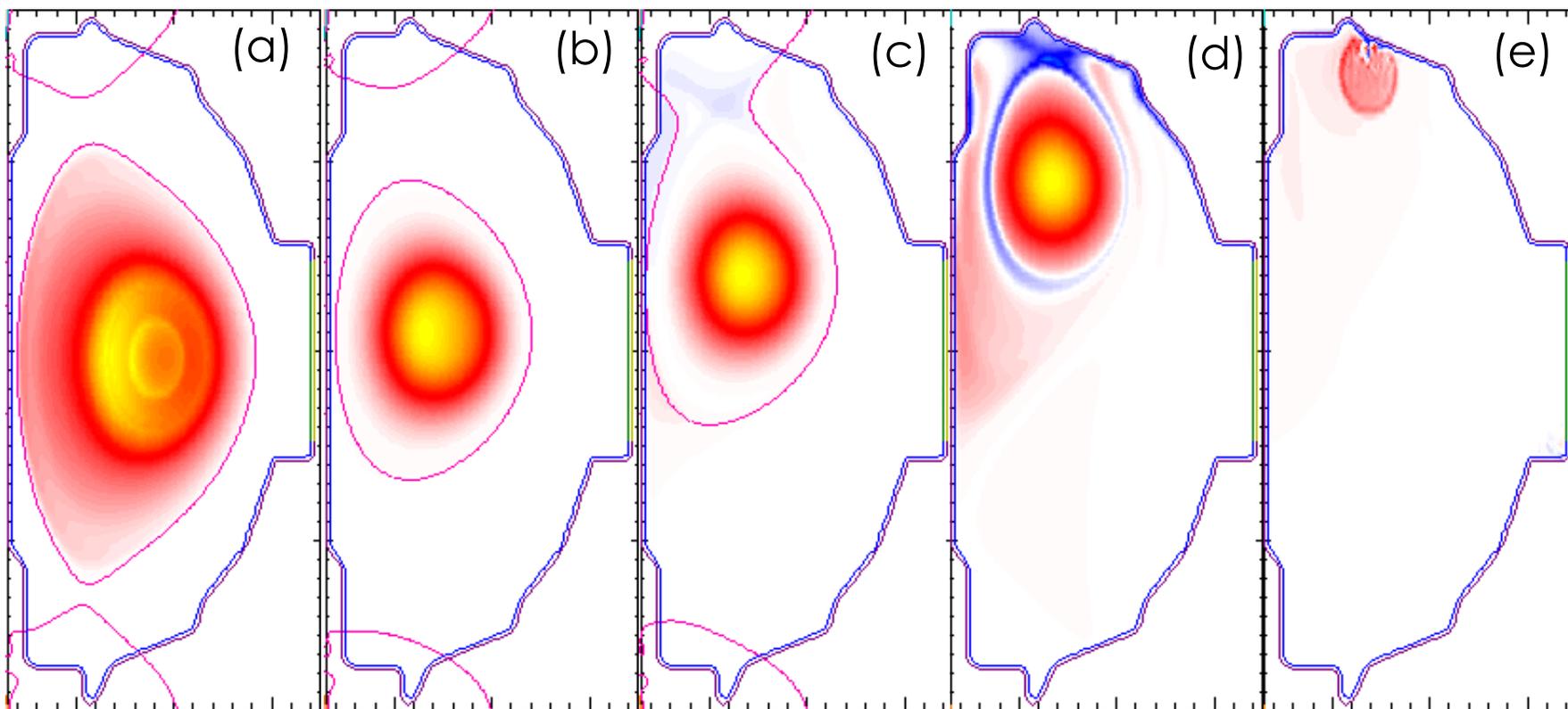
- Maximum J_R occurs very late in VDE, when plasma is limited by lower divertor
- Maximum J_R is roughly 2–2.5 MA/m² in this case
 - Corresponds to $F_Z \sim 500$ kN over ~ 50 cm of the lower divertor



NM Ferraro/TSD/July 2014

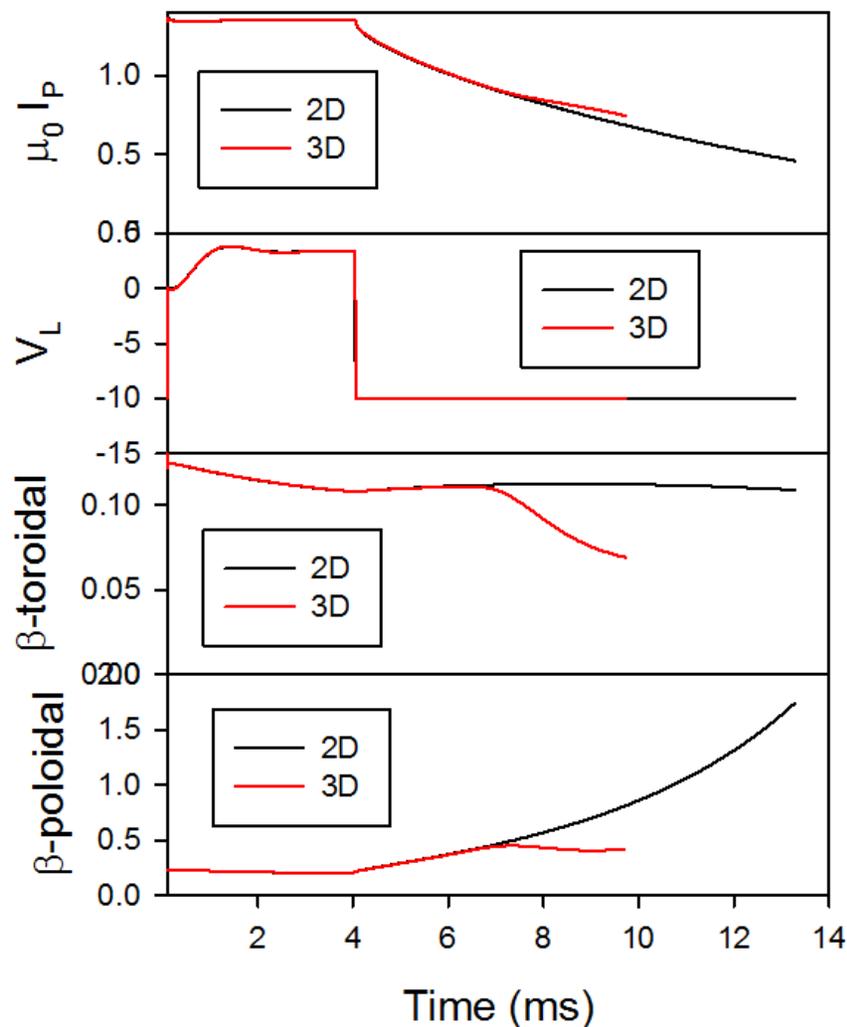
VDEs Also Simulated for NSTX; Results Similar to DIII-D Simulations

- NSTX case also shows **co-current** and **counter-current** current density in the open field-line region



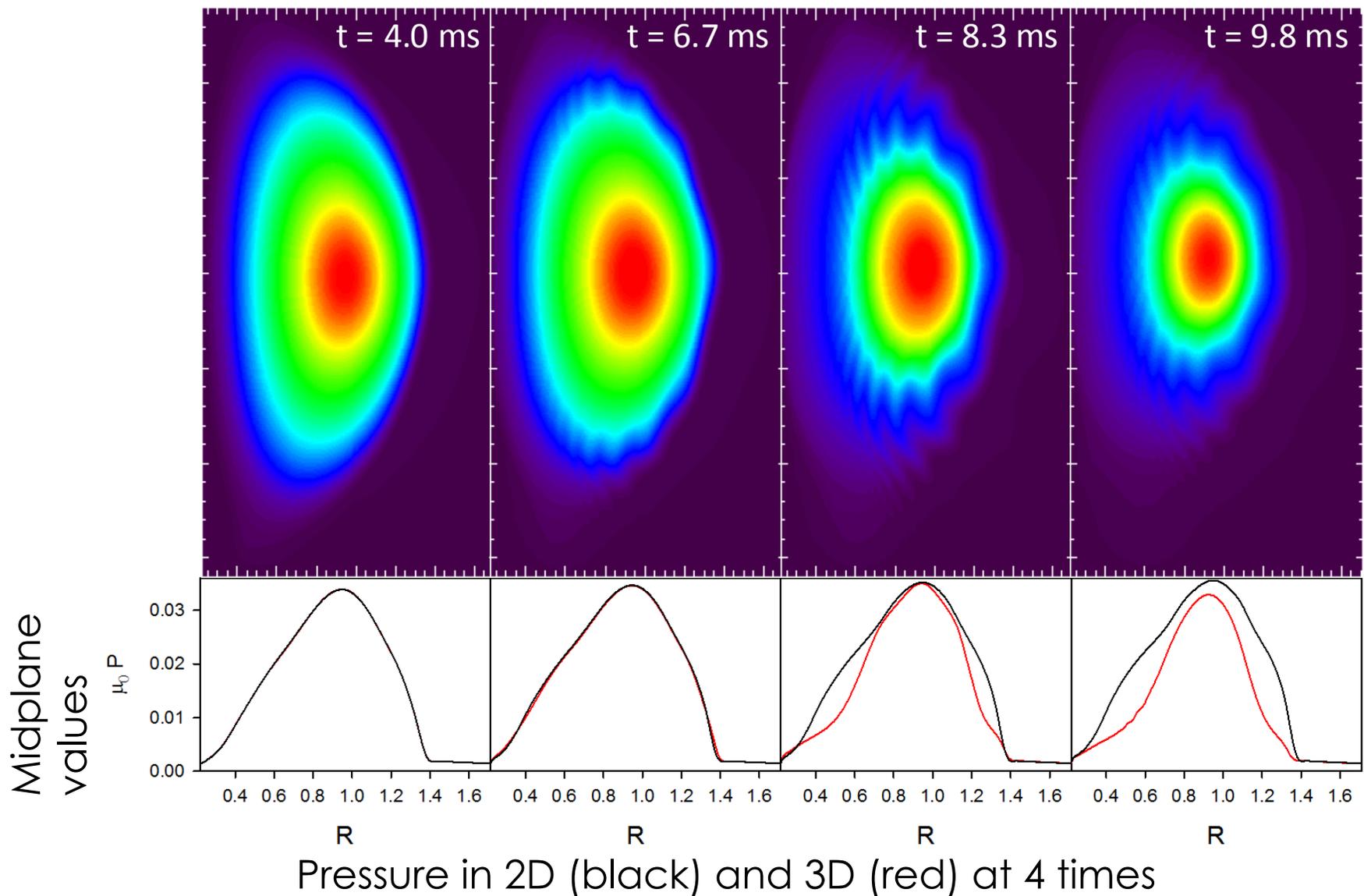
Toroidal current density at 5 times in VDE simulation

Disruptions From V_L Reversal Are Being Explored

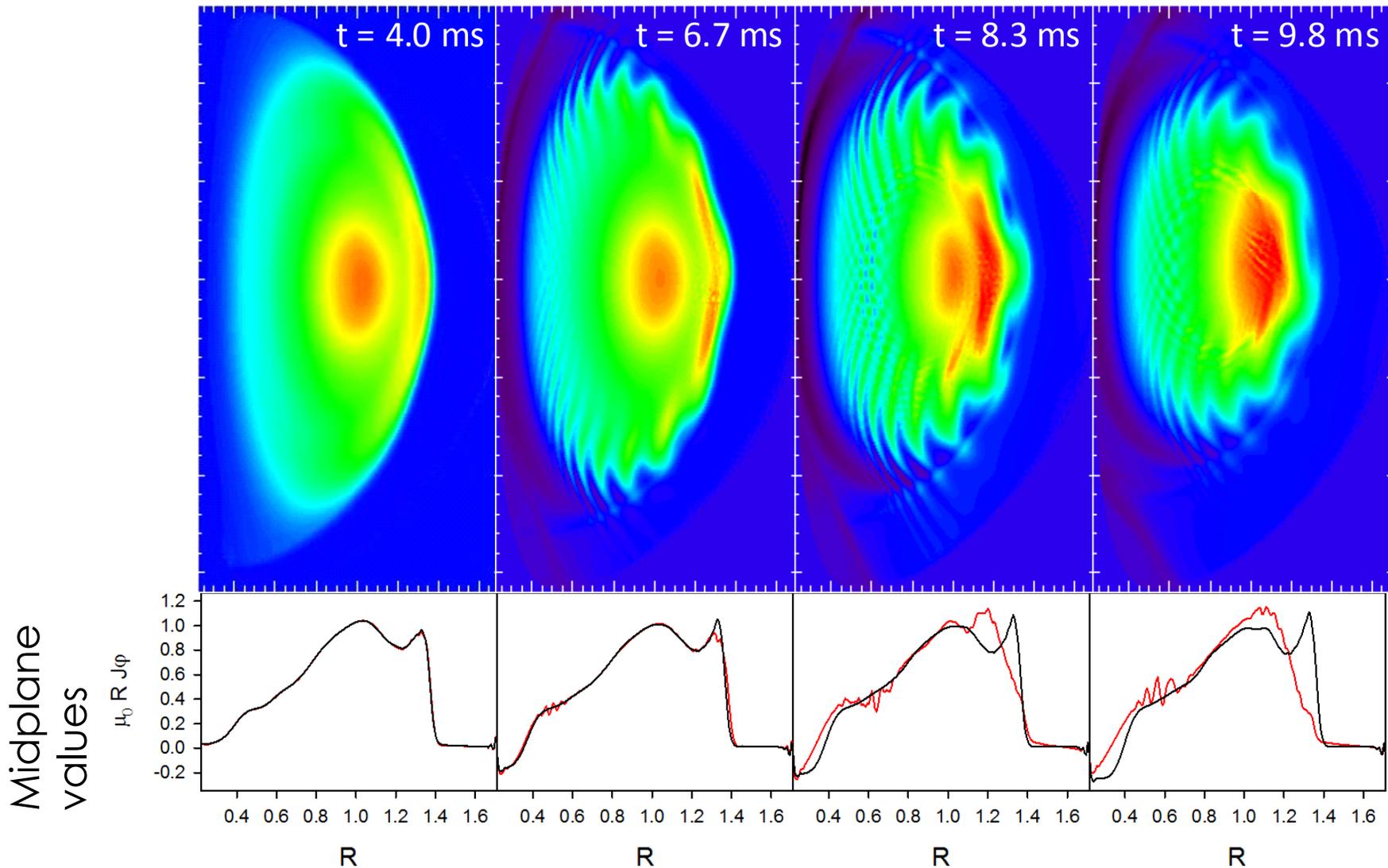


- At the end of the discharge, the loop voltage in NSTX is rapidly reversed to quench the OH coils
- Experimentally, this often leads to a disruption
- We are trying to reproduce this with M3D-C¹ using realistic parameters
- Difference in 2D and 3D behavior is due to 3D instabilities

Edge Mode Found in Reversed- V_L Simulation



Edge Mode Found in Reversed- V_L Simulation



Toroidal current density in 2D (black) and 3D (red) at 4 times

Summary

- **“Free-Boundary” response using resistive wall model is being validated using new 3D magnetic probes**
 - Agreement is encouraging
 - Results sensitive to treatment of open field-line region, especially resistivity; other sensitivities are being explored
- **VDEs successfully simulated in DIII-D and NSTX-U**
 - Axisymmetric, single-fluid
 - Spitzer resistivity, realistic parameters
- **Currents are observed in the wall and open field-line region**
 - At early stages of disruption, wall currents are larger at low η_W
 - At late stages of disruption, wall currents depend weakly on η_W
 - Current spike is observed, and is associated with contraction of plasma and loss of counter- I_p current (not TQ)

Future Work Will Focus on Quantitative Validation and 3D Effects in Disruptions

- **Need cases with lower boundary T_e**
 - Will faster current decay reverse wall current direction?
- **Non axisymmetric instabilities during disruption may lead to sideways forces, enhanced transport**
- **First step: linear stability analysis of 2D evolving equilibrium**
 - Very fast
 - Will show onset, but not saturation / dynamics of $n > 0$ instability
- **Next step: fully nonlinear 3D**
 - This will be possible soon, but expensive
 - Nonlinear evolution will be necessary to quantify forces, etc.

Extra Slides

VDE Calculations Also Successfully Simulated for NSTX

- Position of outer boundary can still strongly affect stability

