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Simulation of Resistive Instabilities in the presence of Runaway Current

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Motivation

During the disruption:

Damage of runaway electrons to the first wall

>During the thermal quench, temperature drops quickly, resistivity increases correspondingly, plasma current starts to decay.

>During the current quench, the current is contributed partly or completely by runaway current.

➢Runaway current profile is usually more peaked than the predisruption current, may induce resistive instabilities.

➢ Resistive instabilities may influence the confinement of runaway electrons.

Motivation

The generation of runaway current (H. Smith, *et.al.*, *PoP*, 2006)
Two of the generation mechanisms: Dreicer and Avalanche

$$\frac{\partial n_{re}}{\partial t} = \frac{\partial n_{re}^{I}}{\partial t} + \frac{\partial n_{re}^{II}}{\partial t}, \qquad (1)$$

Ohm's Law and Faraday's Law

$$J_{''} = \sigma_{''} E_{''} + e n_{re} c, \qquad (2)$$

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial E_{//}}{\partial r}\right) = \mu_0 \frac{\partial J_{//}}{\partial t},\tag{3}$$

Motivation

◆ The initial and final current profile



calculated from the above model



е (тато)

FIG. 2. A simulation of the runaway dynamics in JET discharge 63 133.

H. Smith, et.al., PoP, 2006

Physical Model

Parallel Ohm's Law $E_{II} = \eta (J_{II} - J_{ro})$ (4) Momentum equation $\rho \frac{d\mathbf{v}}{dt} = -\nabla p_c - \nabla \cdot \mathbf{P}_{re} + \mathbf{J} \times \mathbf{B}$ (5)Runaway density equation $\frac{\partial n_{re}}{\partial t} + \nabla \cdot \left(n_{re} v_{//} \mathbf{b} - n_{re} \frac{m_e v_{//}^2}{eB} \mathbf{b} \times \mathbf{\kappa} - n_{re} \frac{\mathbf{b} \times \mathbf{E}}{B} \right) = 0$ (6) $J_{re} = -en_{re}v_{//}, P_{re} = n_{re}m_{e}v_{//}^{2}, |v_{//}| \sim c$ is assumed. Due to this assumption, Eq.(6) is derived from the drift kinetic equation, similar to that of Helander, et. al., PoP, 2007

Implement the above extended-MHD model in the M3D-k code.

Simulation results

- ♦ As a first step, we consider only the effects of runaway current. The effects of runaway pressure tensor are neglected.
- For simplicity, the current is assumed to be completely carried by runaway electrons.

Given the profile of safety factor as



• Two cases: $q_0 = 0.83, 0.93$

◆ Case1: q0=0.83

Without runaway current, the growth rate of 1/1 mode is about 2.14e-2, the mode structures are shown below



• It is the resistive kink mode.

◆ Case 1:q0=0.83

With runaway current, the growth rate of 1/1 mode is about 2.56e-2, the mode structures are shown below



• The growth rate is enhanced slightly, and the mode structures keep almost the same.

The physics: the linearized equations in cylindrical geometry are

$$-(im/q-in)\delta\hat{\phi} + \gamma\delta\psi = \hat{\eta} \left(\nabla_{\perp}^{2}\delta\psi - R\delta\hat{J}_{R}\right),$$
(7)

$$\gamma \nabla^2 {}_{\perp} \delta \hat{\phi} = -\frac{imR}{r} \frac{d\hat{J}_{//0}}{dr} \delta \psi + \frac{in}{q} \left(\frac{m}{n} - q\right) \nabla^2_{\perp} \delta \psi, \qquad (8)$$
$$(m/q - n) \delta J_R - \frac{m}{r} \frac{d\hat{J}_R}{dr} \delta \psi = 0, \qquad (9)$$

where runaway current is a flux function approximately, due to $|v_{ll}| \sim c$. Order analysis: $\delta J_{ll} \sim O(\delta \psi / w_l^2), \delta J_R \sim O(\delta \psi / w_l) \rightarrow \delta J_{ll} > \delta J_R$ The perturbation of runaway current is of higher order. The characteristic functions in the resistive layer without runaway are $\delta \psi = b/2 \Big| xerfc(x/w_l) - (w_l/\sqrt{\pi}) \exp(-(x/w_l)^2) \Big|$ $\delta \phi = i(b/2) \Big(w_l/\sqrt{2} \Big) erfc(x/w_l) \Big|$



The perturbation of runaway current tends to enhance the perturbation of Ohm' current, so that it will increase the growth rate of the resistive kink mode.

With runaway current, without electric drift contribution in runaway equation, the growth rate is 2.45e-2, mode structures are



• The contribution of electric drift in runaway equation is small, due to $c/v_A >> 1$.

With runaway current, without magnetic drift contribution in runaway equation, the growth rate is 2.6e-2, mode structures are



• The effect of magnetic drift velocity is small due to $\Delta_b / \Delta_{layer} >> 1$.

◆ Case2: q0=0.93

Without runaway current, the growth rate of 1/1 mode is about 9.3e-3, the mode structures are



◆ Case 2:q0=0.93

With runaway current, the growth rate of 1/1 mode is about 1.72e-2, the mode structures are



• The growth rate is enhanced, and the mode structures keep almost the same.



The usual sawteeth are observed in the simulation.



Resistive kink mode results a single sawtooth crash only.
A steady state with axi-symmetric equilibrium is reached after the crash.

The plasma current and runaway current become flatten within q=1 surface, q0 greater than 1.

Physics: sawtooth crash tends to flatten current within q=1 surface. After the crash, the runaway current, acting as a current source, stays flatten.

Simulation results

• The evolution of current









Conclusion and discussion

The linear growth rate of 1/1 resistive kink mode is enhanced by the effects of runaway current.

After a single sawtooth crash, plasma reaches a new steady state with axi-symmetric equilibrium.

The profiles of plasma current and runaway current are flatten within q=1 surface.



Thank you for your attention!